

Quantum sensors based on tunneling

From Quantum to Cosmology - 6th June 2024 – Los Alamos



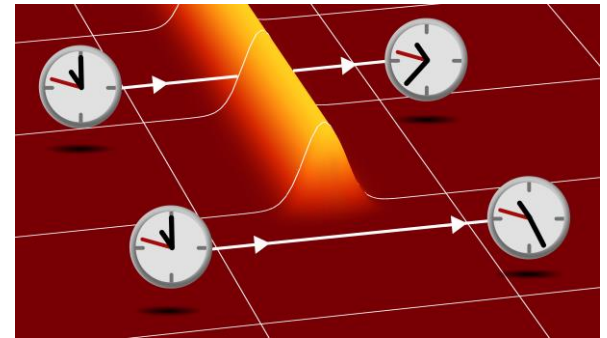
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Theoretical Quantum Optics

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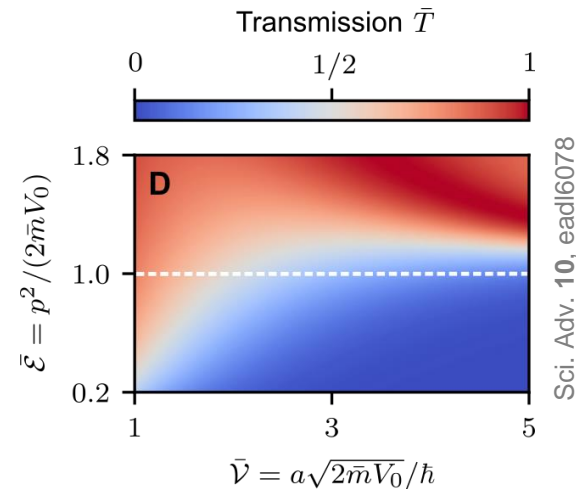
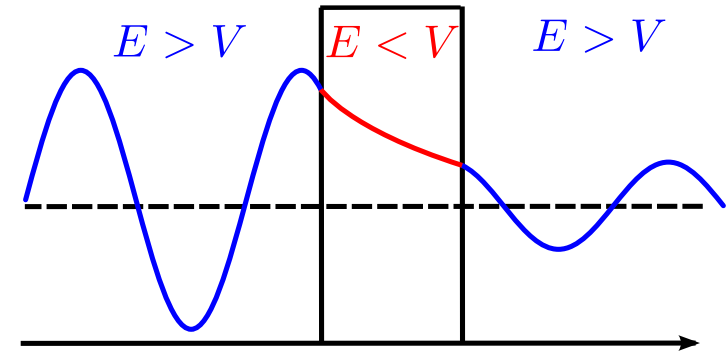
Theoretical quantum
optics group



TECHNISCHE
UNIVERSITÄT
DARMSTADT

Quantum tunneling

- ▶ Transmission of plane waves
- ▶ Analytical solution available
- ▶ Decreasing probability amplitude
- ▶ Exponential decay within barrier region



Quantum tunneling for sensing

Applications

▶ Nuclear fusion

Rev. Mod. Phys. **70**, 77

▶ Josephson junctions

Phys. Rev. Lett. **95**, 010402

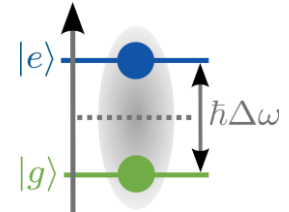
▶ Scanning tunneling microscopy

▶ Hawking radiation

J. High Energy Phys. JHE09(2005)037

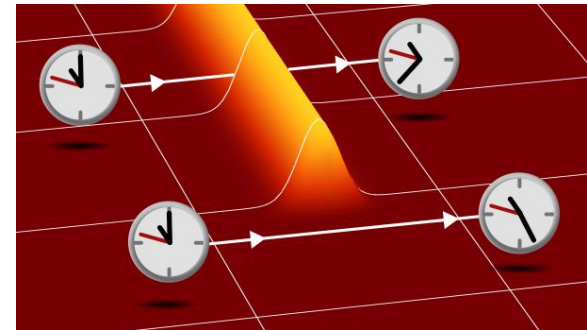
Quantum clocks

▶ Monitor transition frequency



▶ Operational approach

Time is what you read off a clock



Inertial sensing

Atom interferometer

- ▶ Gyroscope
 - ▶ Sagnac effect

Phys. Rev. A **80**, 063604

- ▶ Gravimeter

- ▶ Linear acceleration

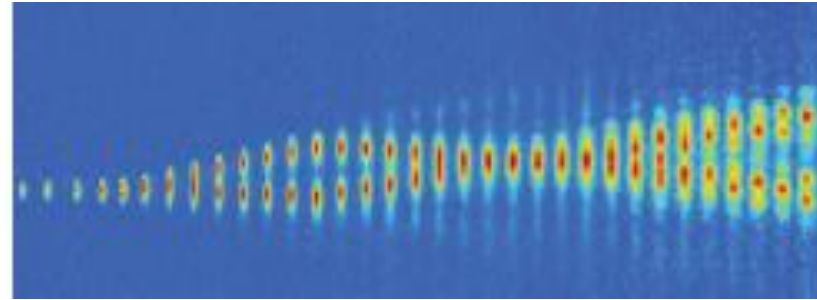
Phys. Rev. Lett. **117**, 203003

- ▶ Gradiometer

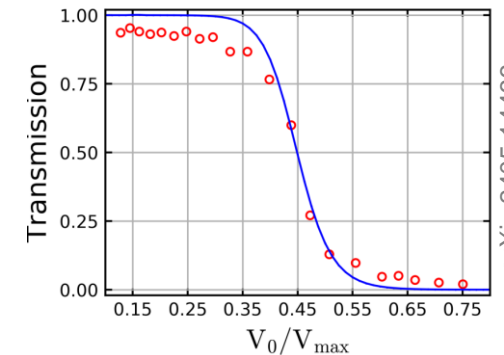
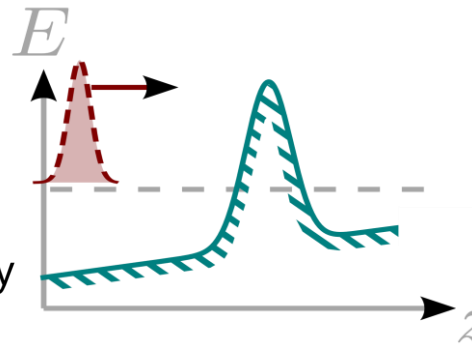
- ▶ Curvature of gravity

Science **315**, 74

Idea: Quantum Tunneling for gravimetry

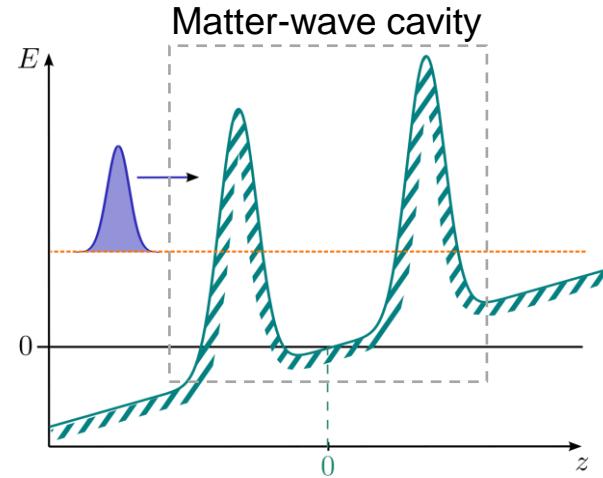
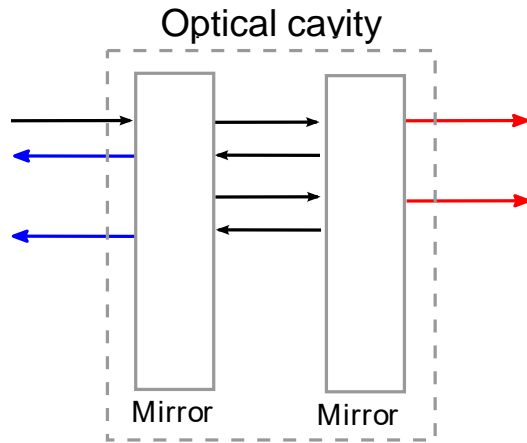


Phys. Rev. Lett. **110**, 093602



arXiv:2405.14466

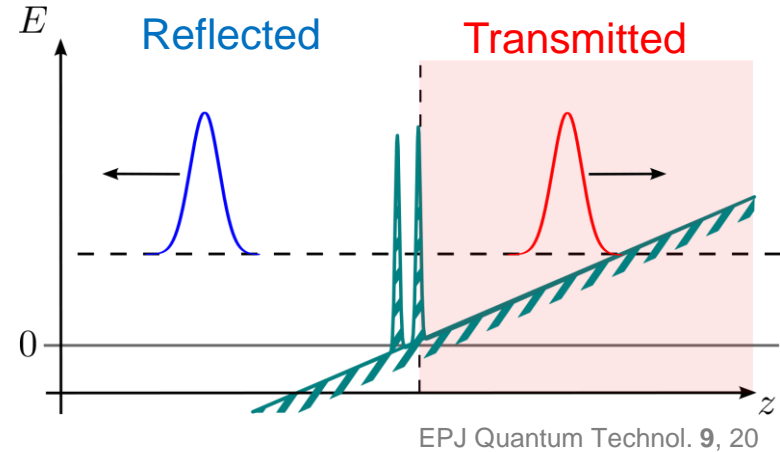
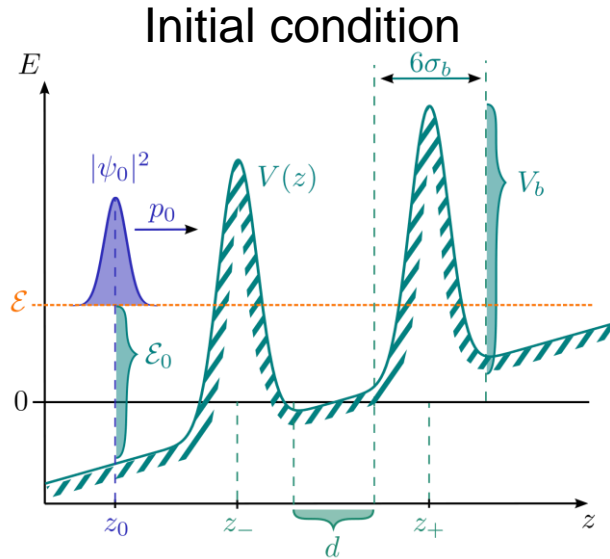
Matter-wave Fabry-Pérot interferometer



EPJ Quantum Technol. 9, 20



Matter-wave Fabry-Pérot interferometer



- ▶ Center of cavity as reference for kinetic energy $\mathcal{E} = \mathcal{E}_0 - mg|z_0|$
- ▶ Accounts for effects of gravity prior to scattering

Transmission spectrum

Transmission spectrum for plane waves

Sci. Rep. 10, 15052

Momentum representation

$$T_R = \int_{-\infty}^{\infty} dp \, T(p) |\psi_0(p)|^2$$

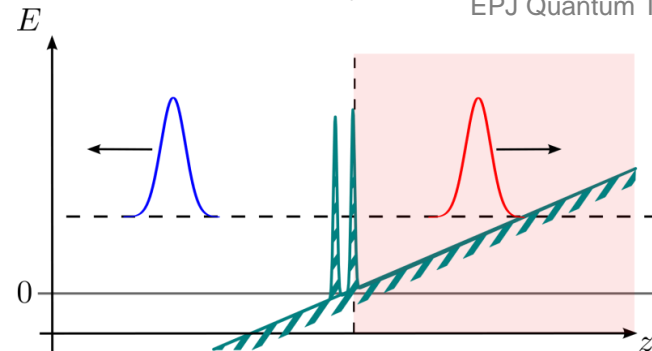
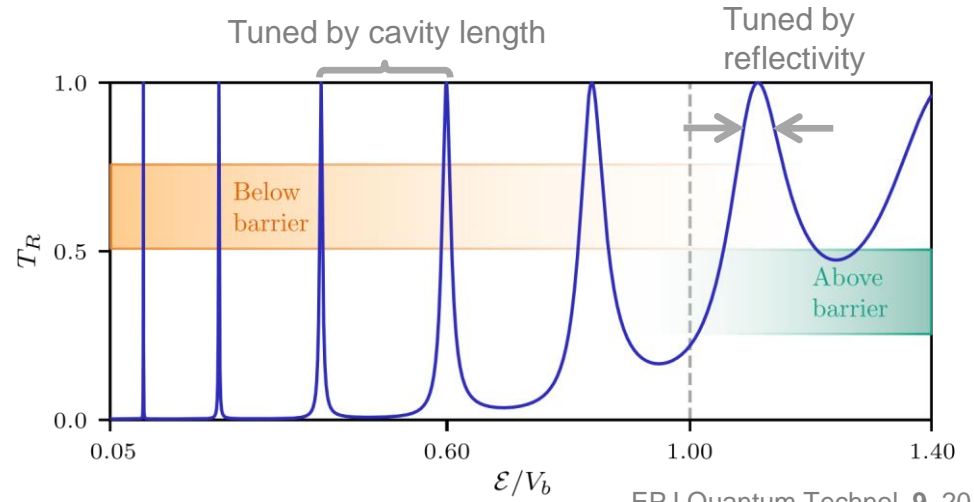
↑ Transmission amplitudes ↑ Initial wave packet

→ Momentum filtering

Position representation

$$T_R = \int_{z_+}^{\infty} dz \, |\psi_f(z)|^2$$

↑ Scattered wave packet



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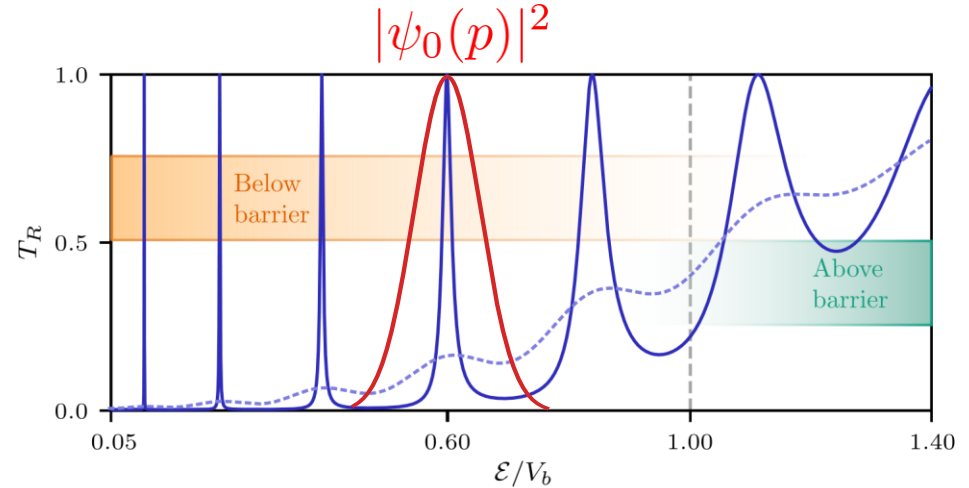
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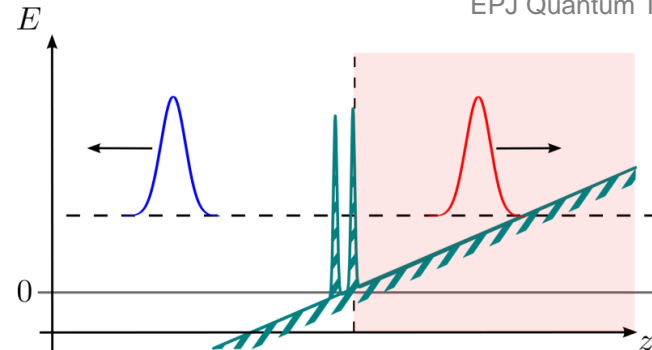
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EPJ Quantum Technol. 9, 20



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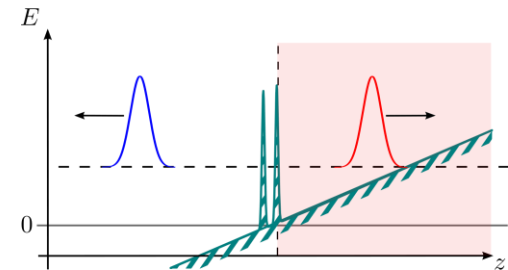
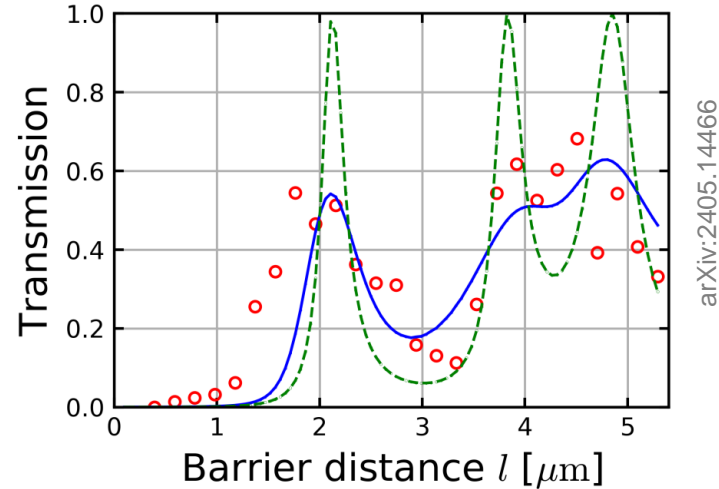
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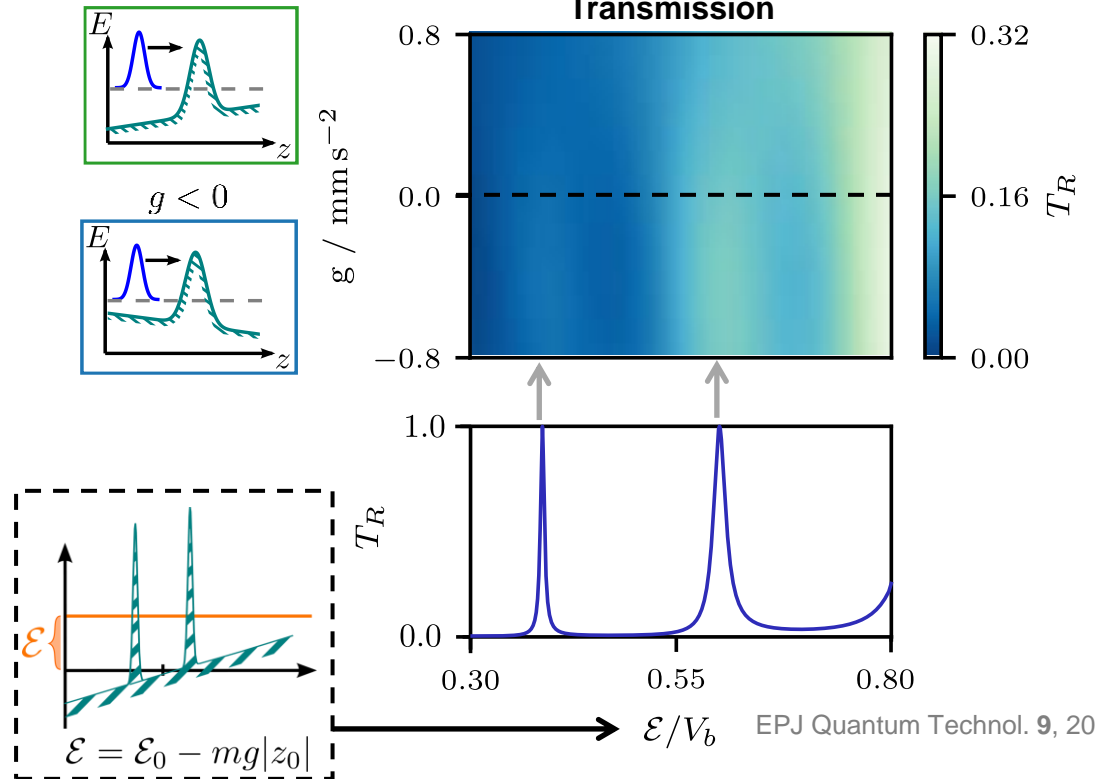


Transmission of a wave packet

Observable:

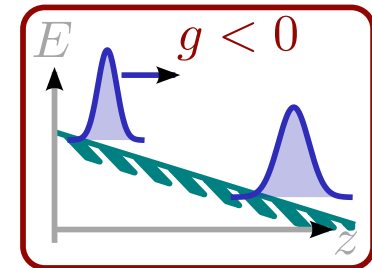
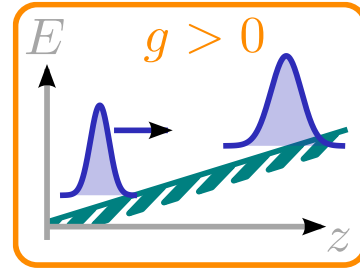
Fraction of transmitted atoms

- ▶ Resonances depend on gravity
- ▶ Main effects prior to scattering
 - 1) Acceleration of wave packet
→ Shift of resonances
 - 2) Wave packet deformation



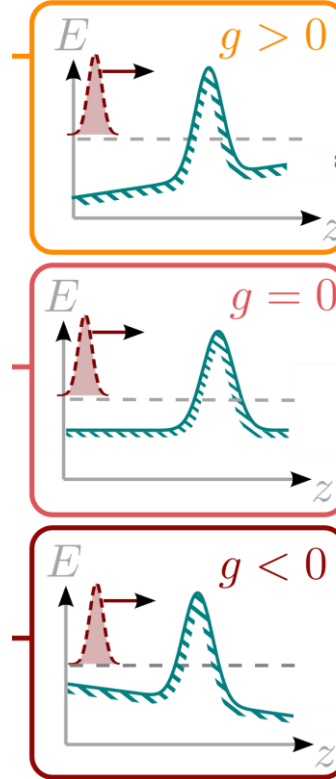
Wave packet deformation

- ▶ Linear potential $V(z) = mgz$
 - ➔ No effect on width of wave packet



Wave packet deformation

- ▶ Linear potential $V(z) = mgz$
 - ➔ No effect on width of wave packet
- ▶ Gaussian barriers $V(z) = mgz + V_{\text{Gauss}}$
 - ➔ Deformation of wave packet



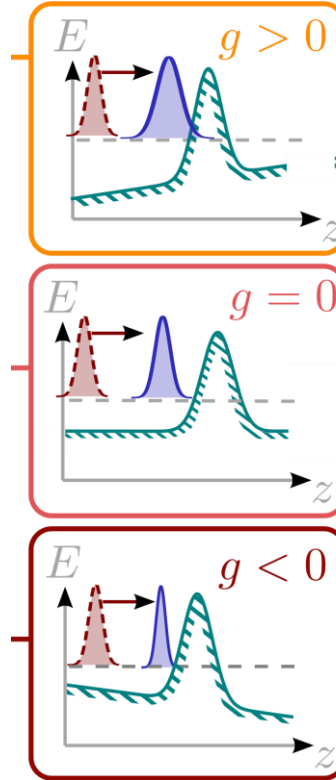
EPJ Quantum Technol. 9, 20

Wave packet deformation

- ▶ Linear potential $V(z) = mgz$
 - ➔ No effect on width of wave packet

- ▶ Gaussian barriers $V(z) = mgz + V_{\text{Gauss}}$
 - ➔ Deformation of wave packet

- ▶ Prominence of resonances modified



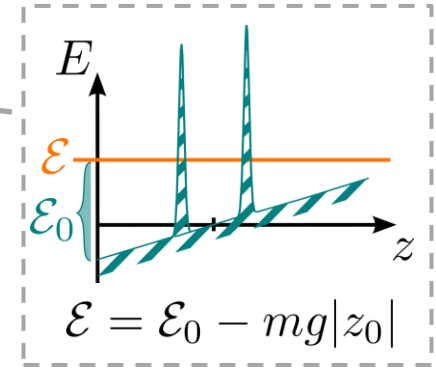
EPJ Quantum Technol. 9, 20

Sensitivity of a wave packet

Gaussian error propagation $T_R = T_R(g, \mathcal{E})$

$$\Delta T_R^2 = |\partial_g T_R|^2 \Delta g_R^2 + |\partial_{\mathcal{E}} T_R|^2 \Delta \mathcal{E}^2$$

$$\Delta \mathcal{E} = |m z_0| \Delta g_R$$



Relative uncertainty

$$\delta g_R = \frac{\Delta g_R}{\sqrt{\nu} \sqrt{N} g} = \frac{1}{\sqrt{\nu} \sqrt{N}} \frac{\Delta T_R}{\sqrt{|g \partial_g T_R|^2 + |m g z_0 \partial_{\mathcal{E}} T_R|^2}}$$

Number of
measurements

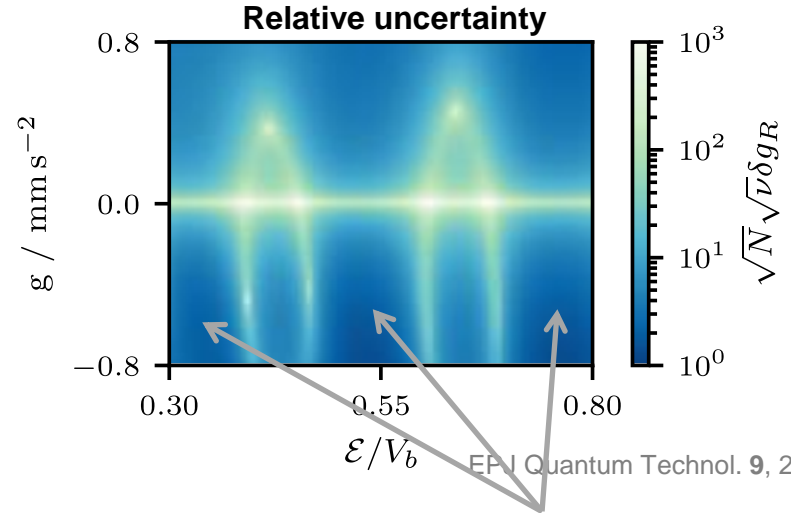
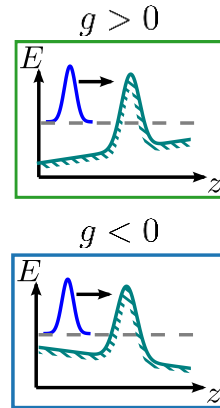
Particle
number

Sensitivity of
transmission
coefficient

Propagation prior to
impact

Sensitivity of a wave packet

Observable: Fraction of transmitted atoms



Relative uncertainty

$$\delta g_R = \frac{\Delta g_R}{\sqrt{\nu} \sqrt{N} g} = \frac{1}{\sqrt{\nu} \sqrt{N}} \frac{\Delta T_R}{\sqrt{|g \partial_g T_R|^2 + |m g z_0 \partial_E T_R|^2}}$$

Number of measurements $\rightarrow \sqrt{\nu}$ Particle number $\rightarrow \sqrt{N}$

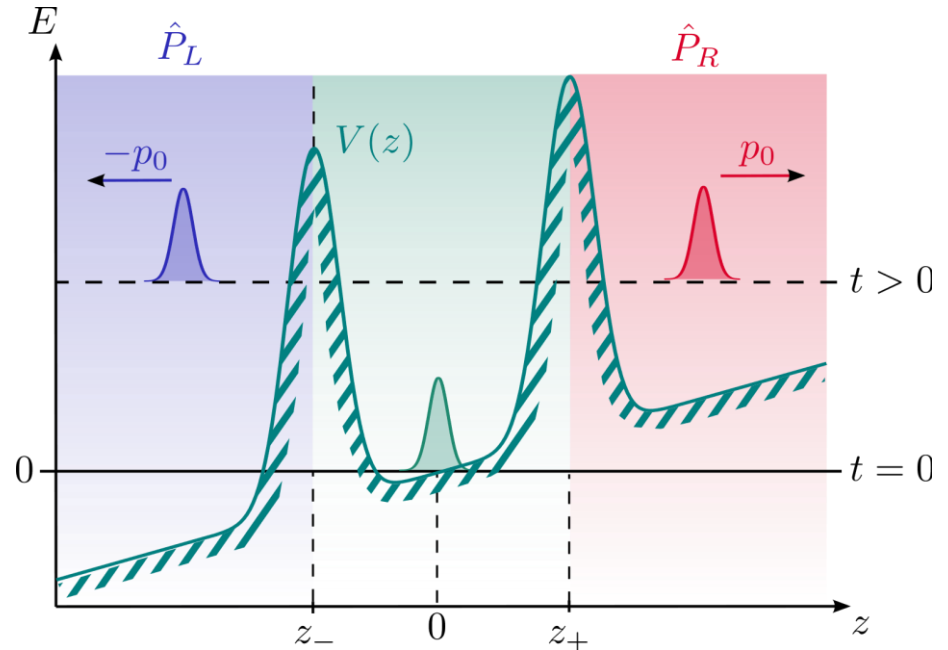
Sensitivity of transmission coefficient $\rightarrow |g \partial_g T_R|^2$ Propagation prior to impact $\rightarrow |m g z_0 \partial_E T_R|^2$

Working points

Asymmetric transmission

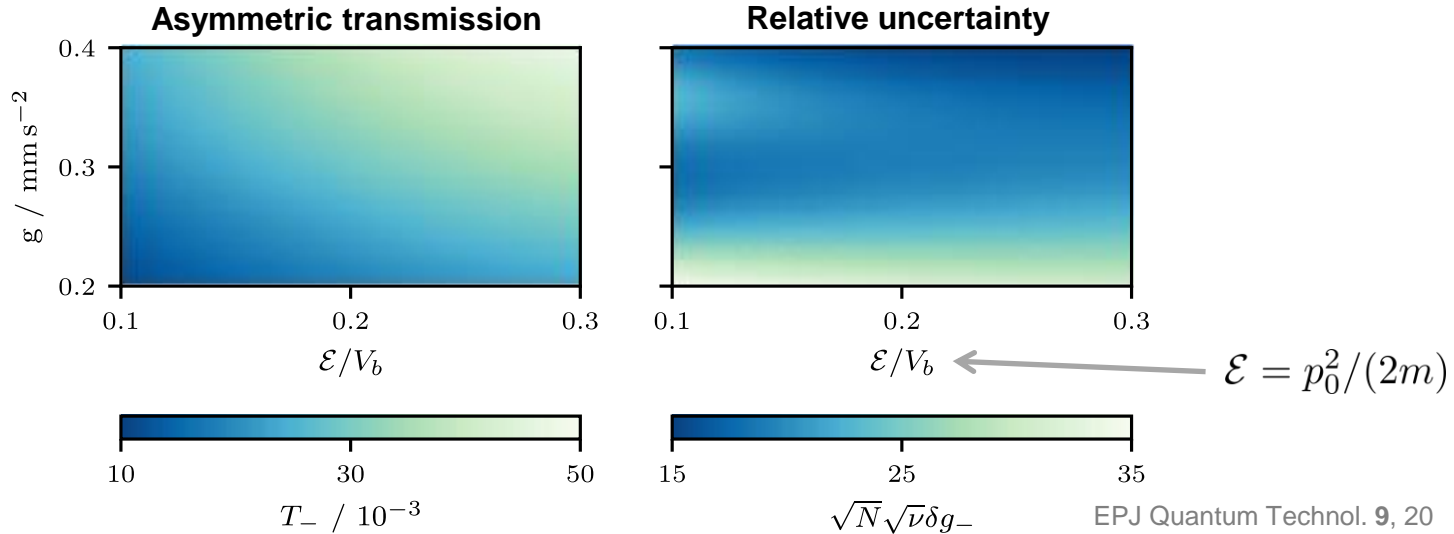
- ▶ Initial state: Kicked wave packet
- ▶ Momentum transfer via Double Bragg
Phys. Rev. A **88**, 053608
- ▶ No acceleration
 → Symmetric motion
- ▶ Observable: Asymmetry

$$T_- = T_L - T_R$$



EPJ Quantum Technol. **9**, 20

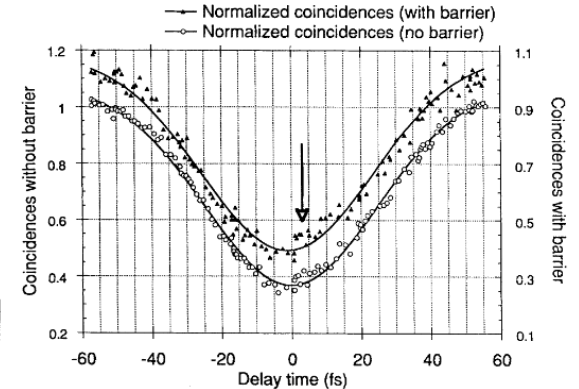
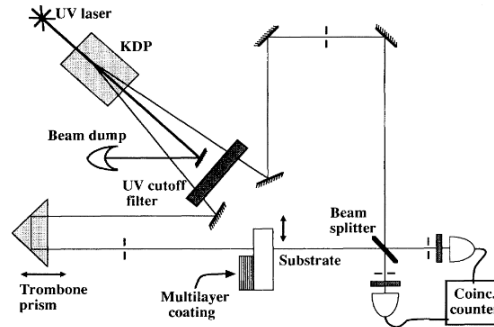
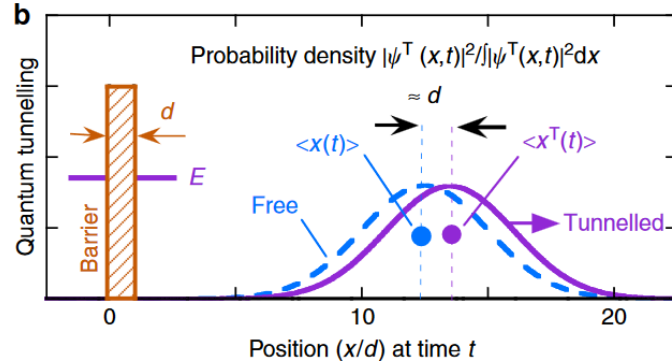
Asymmetric transmission



- ▶ No structure of transmission spectrum visible
- ▶ Sensitivity one order of magnitude smaller than for T_R measurement

Concept on tunneling times

- ▶ Relativity: Proper time measured by an ideal clock along a worldline
- ▶ Quantum tunneling
 - ▶ Impossible to assign a worldline
- ▶ Arrival time
 - Phys. Rev. Lett. **49**, 1739
 - Rev. Mod. Phys. **61**, 917
 - Commun. Phys. **1**, 47
- ▶ Shift in position upon arrival
- ▶ Hong-Ou-Mandel setup
 - Phys. Rev. Lett. **71**, 708
- ▶ Attoclock experiments
 - Science **322**, 1525
 - Nature **446**, 627



Concept on tunneling times

► Interaction (dwell) time

Phys. Rev. A **107**, 042216

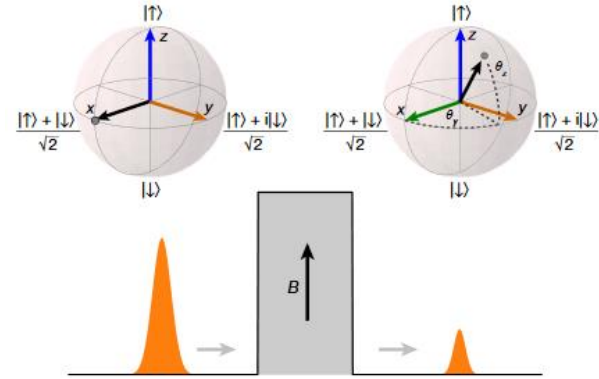
Phys. Rev. B **27**, 6178

Nature **583**, 529

► Time spent in barrier

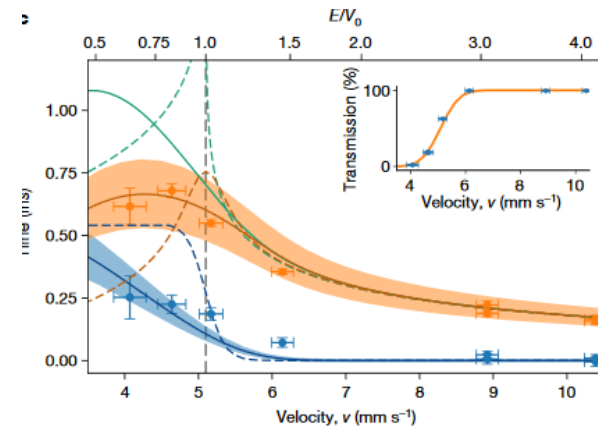
► Larmor clock Phys. Rev. Lett. **127**, 133001

► Phase shift acquired while in barrier



Nature **583**, 529

Phys. Rev. Lett. **71**, 708



Nature **583**, 529

Time is what you read off a clock



Atom with internal
structure

Relativistic effects

- ▶ Atom interferometers for relativistic physics

- ▶ Violations of equivalence principle

Phys. Rev. D **107**, 064007

- ▶ Gravitational waves

Phys. Rev. Lett. **110**, 171102

- ▶ Dark matter candidates

AVS Quantum Sci. **5**, 044404

- ▶ Mass defect $E = mc^2$

- ▶ Different internal states have different mass

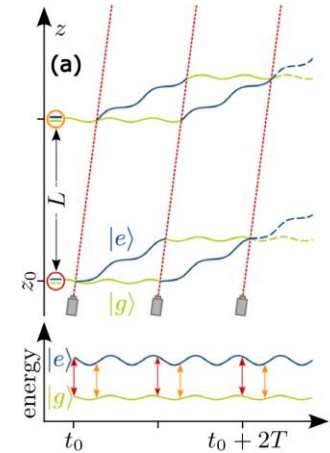
PRX Quantum **5**, 020322

Phys. Rev. A **98**, 042106

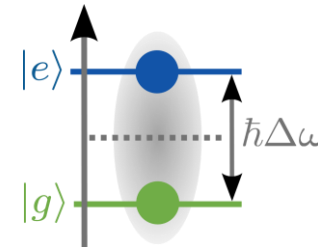
Laser Phys. Lett. **15**, 035703

Phys. Rev. A **100**, 052116

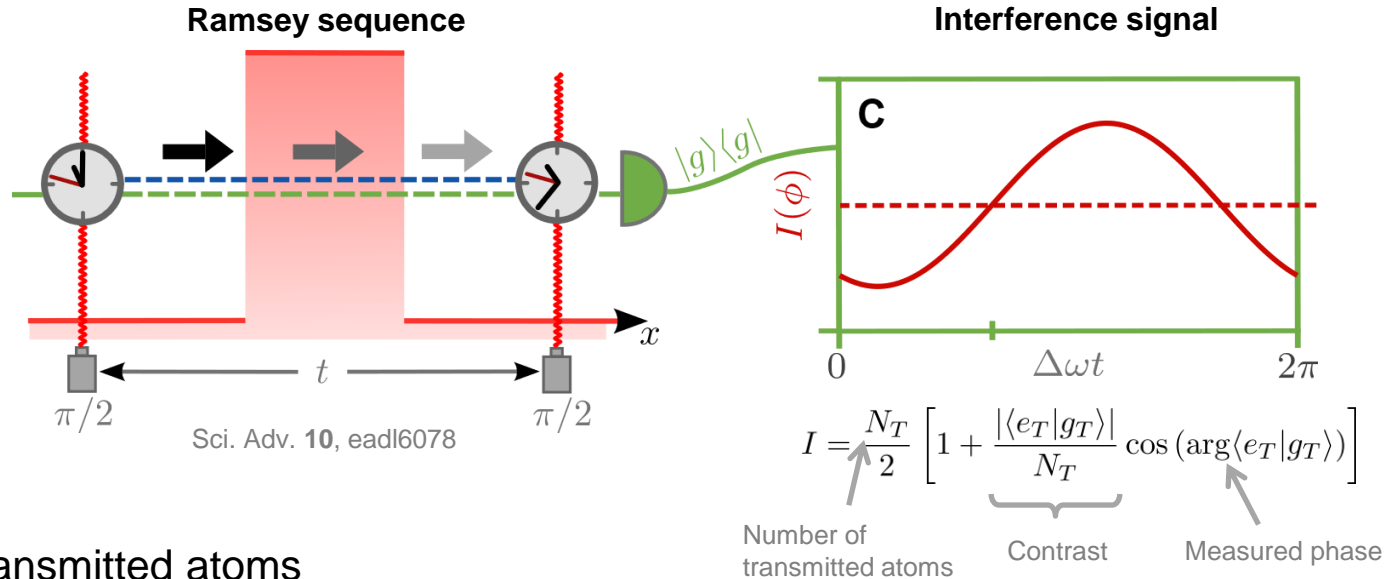
$$m_{e/g} = \bar{m} \pm \frac{\Delta m}{2} = \bar{m} \pm \frac{\hbar \Delta \omega}{2c^2} \leftarrow \text{clock frequency}$$



AVS Quantum Sci. **5**, 044404



Ramsey Clock



State of transmitted atoms

$$|e_T/g_T\rangle = t_{e/g}(E) \exp \left[-i \left(\omega_{e/g} t - \frac{p^2}{2m_{e/g} \hbar} t \right) \right] |e_{in}/g_{in}\rangle$$

Transmission coefficient

$$|t_{e/g}| e^{i\varphi_{e/g}} \cong \sqrt{T} e^{i\bar{\varphi} \pm i\Delta\omega t}$$

Phase contribution

Phase difference measured by a tunneled clock

$$|e_T/g_T\rangle = t_{e/g}(E) \exp \left[-i \left(\omega_{e/g} t - \frac{p^2}{2m_{e/g} \hbar} t \right) \right] |e_{in}/g_{in}\rangle$$

$$\arg\langle e_T|g_T\rangle = \Delta\omega(t + \delta t + \tau) - [\phi(t) - \phi(0)]$$

e.g. laser phase

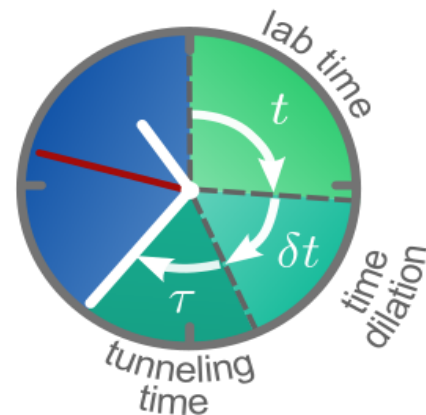
Lab time Time dilation Tunneling time

$$\delta t = \frac{1}{2} \left(\frac{p}{\bar{m}c} \right)^2 t$$

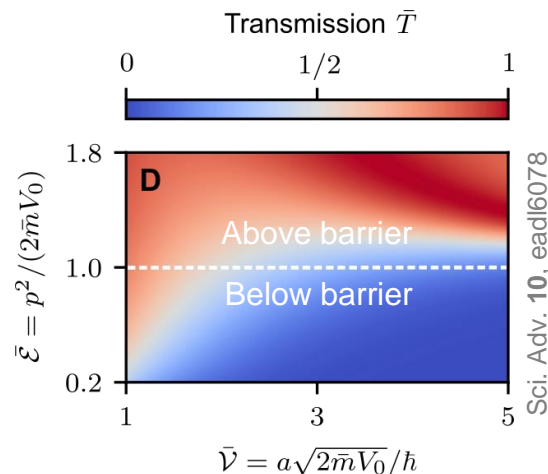
Tunneling time

$$\tau = \frac{\bar{V}\bar{T}}{4\bar{\omega}\sqrt{\bar{\epsilon}}(\bar{\epsilon}-1)} \left[(2\bar{\epsilon}-1) - \frac{\sinh(2\bar{V}\sqrt{1-\bar{\epsilon}})}{2\bar{V}\sqrt{1-\bar{\epsilon}}} \right]$$

\uparrow incoming energy $\bar{\epsilon} = \frac{p^2}{2\bar{m}V_0}$ \uparrow barrier parameter $\bar{V} = a\sqrt{2\bar{m}V_0}/\hbar$



Sci. Adv. 10, eadl6078



Sci. Adv. 10, eadl6078

Tunneling time

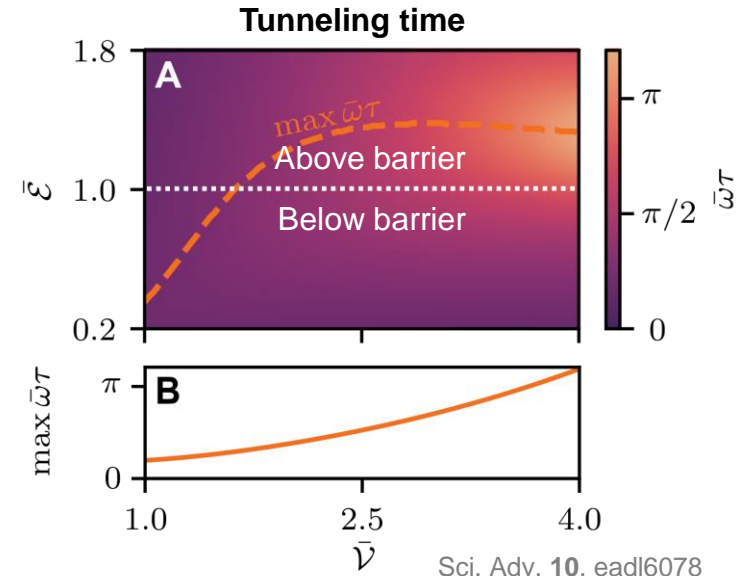
Tunneling time

$$\tau = \frac{\bar{V}T}{4\bar{\omega}\sqrt{\bar{\epsilon}}(\bar{\epsilon} - 1)} \left[(2\bar{\epsilon} - 1) - \frac{\sinh(2\bar{V}\sqrt{1 - \bar{\epsilon}})}{2\bar{V}\sqrt{1 - \bar{\epsilon}}} \right]$$

↑ incoming energy $\bar{\epsilon} = \frac{p^2}{2mV_0}$
↑ barrier parameter $\bar{V} = a \frac{\sqrt{2mV_0}}{\hbar}$

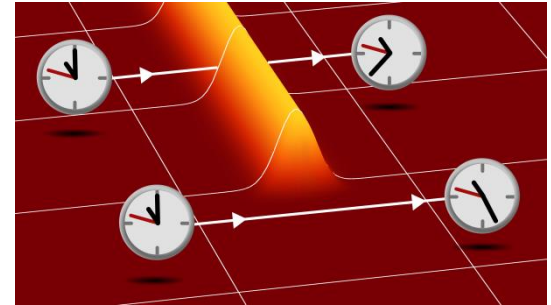
- ▶ Short tunneling times for
 - ▶ small tunneling probabilities
 - ▶ far above barrier

- ▶ No clear distinction between tunneling and classical regime



Differential measurement

- ▶ Differential measurement scheme to isolate tunneling phase
- ▶ Initial state: Cold atomic cloud
- ▶ Measurement of ground-state population

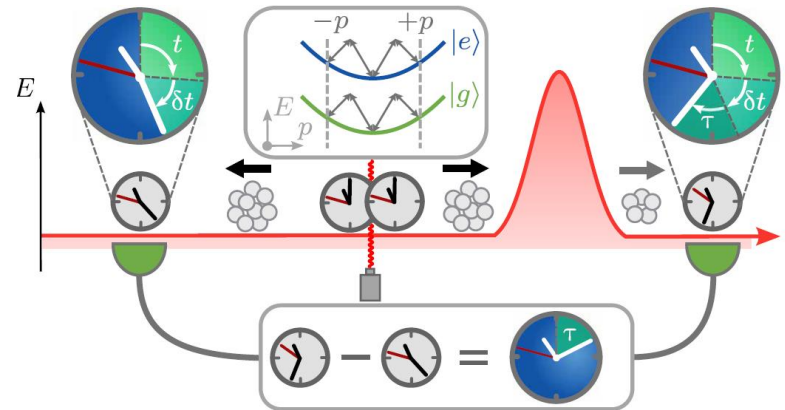


Time dilation

$$\delta t = \frac{t}{2N_T} \int dp \bar{T}(p) |\psi_0(p)|^2 \left(\frac{p}{\bar{m}c} \right)^2$$

Tunneling time

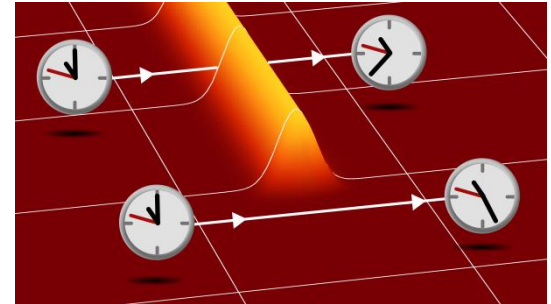
$$\tau = \frac{1}{N_T} \int dp \bar{T}(p) |\psi_0(p)|^2 \tau(p)$$



Sci. Adv. **10**, ead16078

Differential measurement

- ▶ Differential measurement scheme to isolate tunneling phase
- ▶ Insensitive to position uncertainties
- ▶ Velocity uncertainty \rightarrow additional contributions



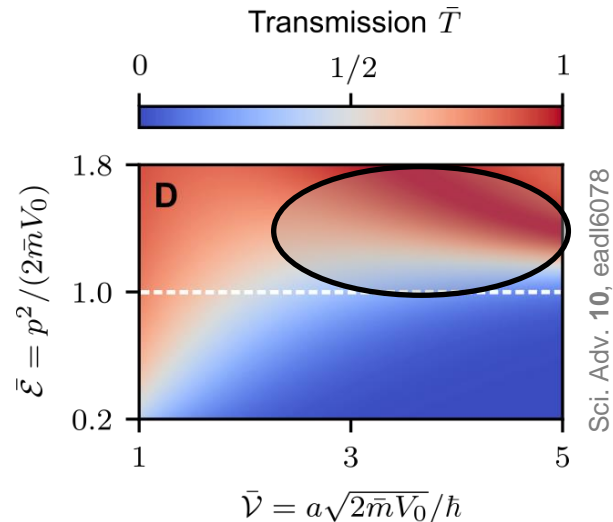
- ▶ But for $|\partial_p \bar{T}|_{p_0} \delta p \ll 1$

$\bar{T}(p) \approx 1$ \nearrow Momentum width
 \nearrow \rightarrow Delta-kick collimation

Phys. Rev. Lett. **78**, 2088

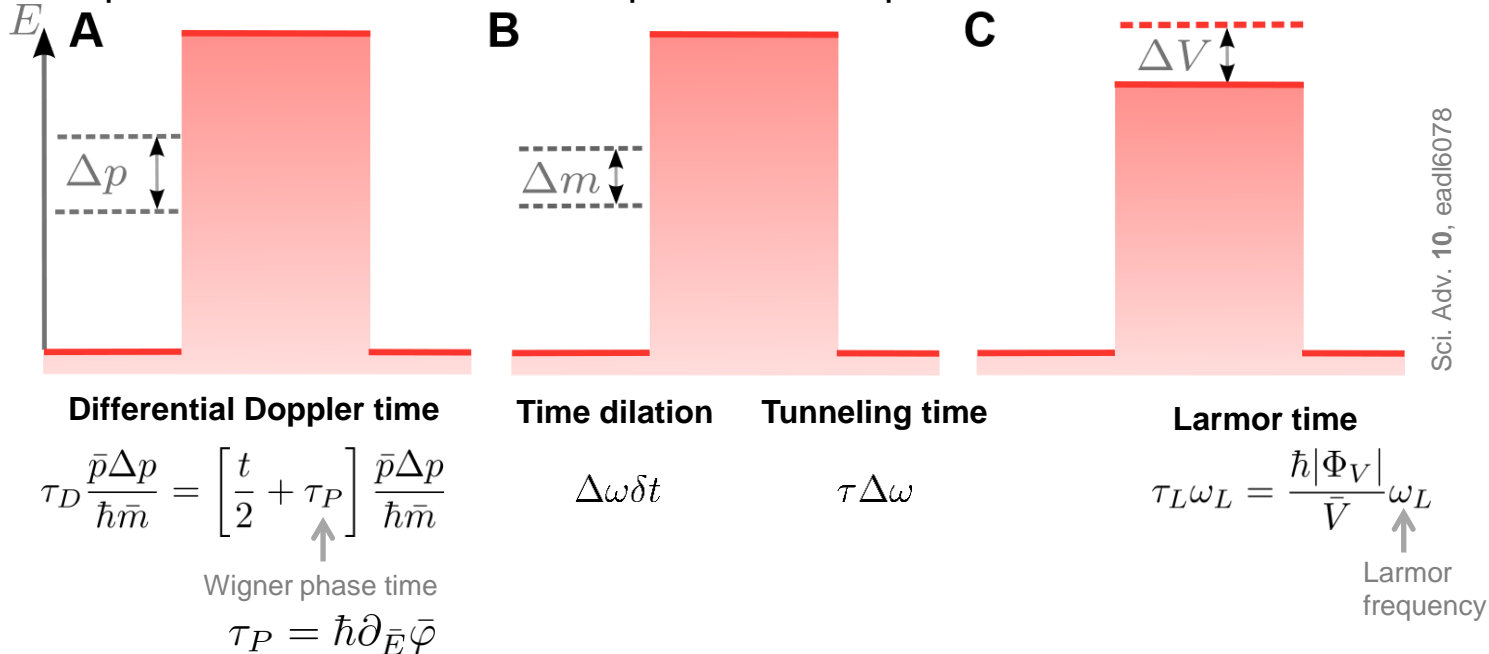
Time dilation

$$\delta t = \frac{t}{2\hbar \cancel{V}} \int dp \cancel{T}(p) |\psi_0(p)|^2 \left(\frac{p}{\bar{m}c} \right)^2$$



Experimental imperfections

- Additional phase contributions from experimental imperfections

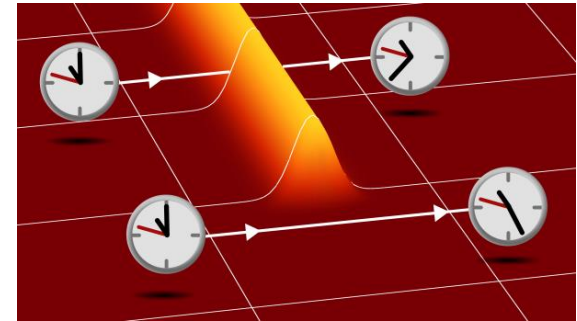
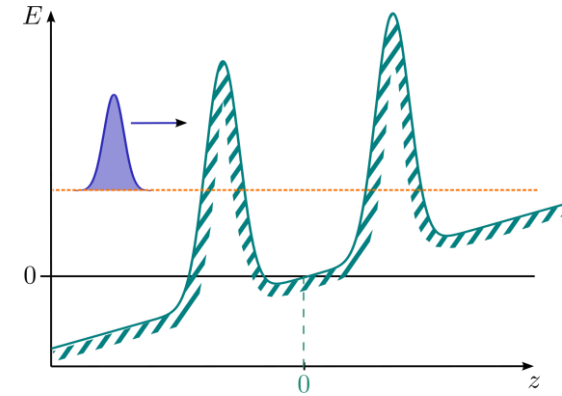


Sci. Adv. 10, ead6078

Phys. Rev. 98, 145

Conclusion

- ▶ Experiments currently performed
- ▶ „Time is what you read off a clock“
 - ▶ No superluminal or instantaneous tunneling
- ▶ Unification of definition of tunnelling times



Thank you for your attention



P. Schach, A. Friedrich, J. R. Williams, W. P. Schleich & E. Giese

Tunneling gravimetry

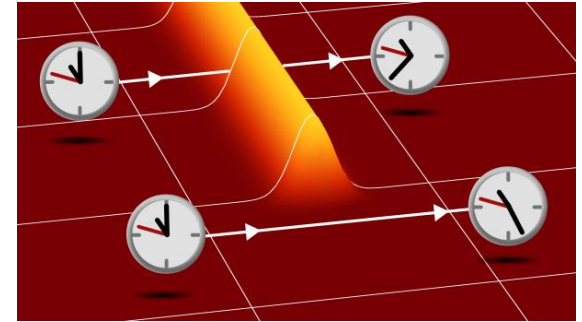
EPJ Quantum Technology **9**, 20 (2022)



P. Schach & E. Giese

A unified theory of tunneling times promoted by Ramsey clocks

Science Advances **10**, eadl6078 (2024)



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